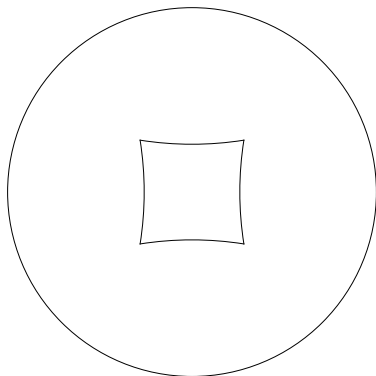


Do we live in a hologram ?

Antoine Bourget
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Rencontres des Jeunes Physicie(n)n(e)s

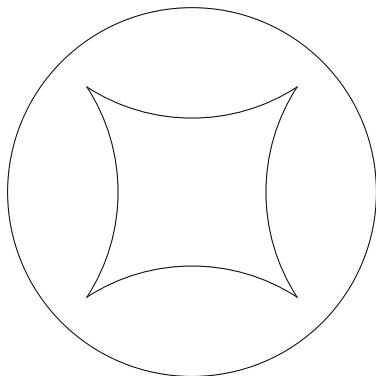
November 10, 2015

Hyperbolic geometry : the Poincaré Disk



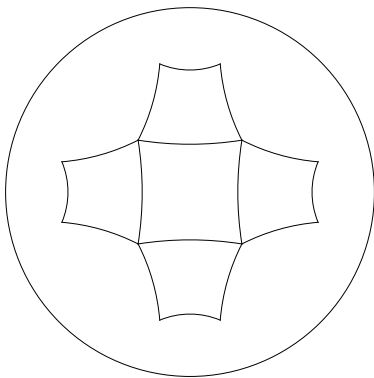
$$ds^2 = 4 \frac{dr^2 + r^2 d\theta^2}{(1 - r^2)^2}$$

Hyperbolic geometry : the Poincaré Disk



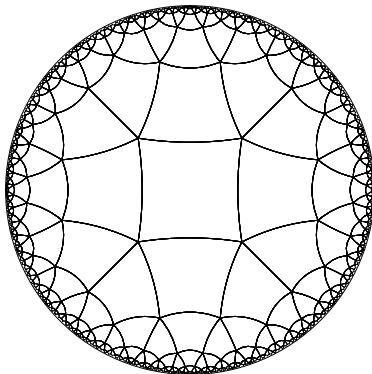
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Hyperbolic geometry : the Poincaré Disk



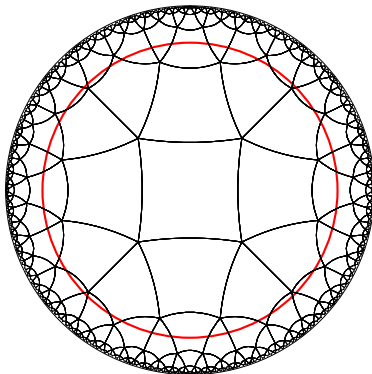
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Hyperbolic geometry : the Poincaré Disk



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Hyperbolic geometry : the Poincaré Disk



$$\text{Radius} = R = 2 \operatorname{argth} r \quad (1)$$

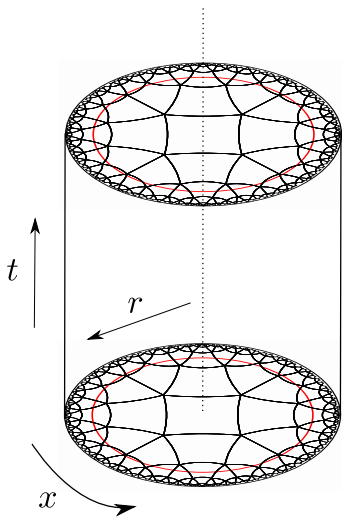
$$\text{Perimeter} = 2\pi \sinh R \quad \text{Area} = 2\pi(\cosh R - 1) \quad (2)$$

The Holographic Principle

In hyperbolic space-times (Anti-de Sitter), we could choose to describe reality with two different sets of laws :

- 3 space dimensions with gravity
- 2 space dimensions without gravity

This would give a complete and logically consistent quantum description of gravity.



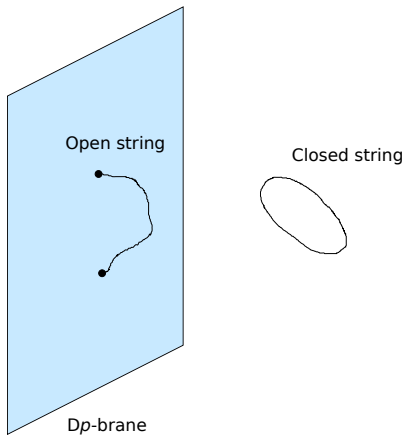
Do we live in a hologram ?

- A crucial aspect of AdS spacetime is that it has a boundary where time is well defined. In an expanding universe which started with a Big Bang, where can we put the boundary theory ?
- What is the spacial curvature of our universe ?
- Mathematical difficulties

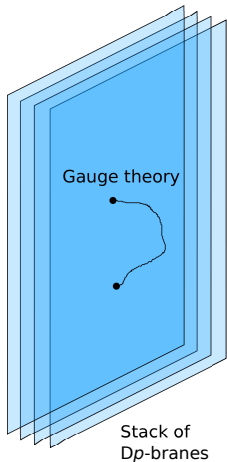
But many other applications !

- Understand better the general properties of gravity
- Shed new light on the physics of strongly coupled systems, for instance in Condensed Matter Theory
- Description of quark-gluon plasma at high energy

D-branes in string theory



D-branes and gauge theories



On the worldvolume of a stack of N D3-branes lives a 4-dimensional $SU(N)$ gauge theory at low energy.

The coupling between closed and open strings is controlled by $\lambda = g_s N \ll 1$.
An observer sees two decoupled theories :

- A $SU(N)$ gauge theory from the open strings;
- A (super)gravity theory in flat spacetime from closed strings.

Black branes

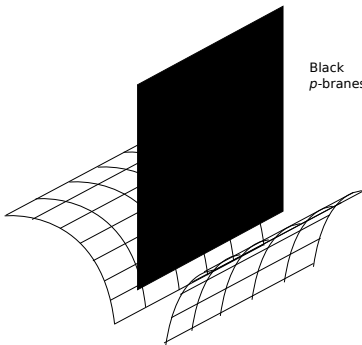
The brane is now very massive and curves spacetime:

Black
 p -branes

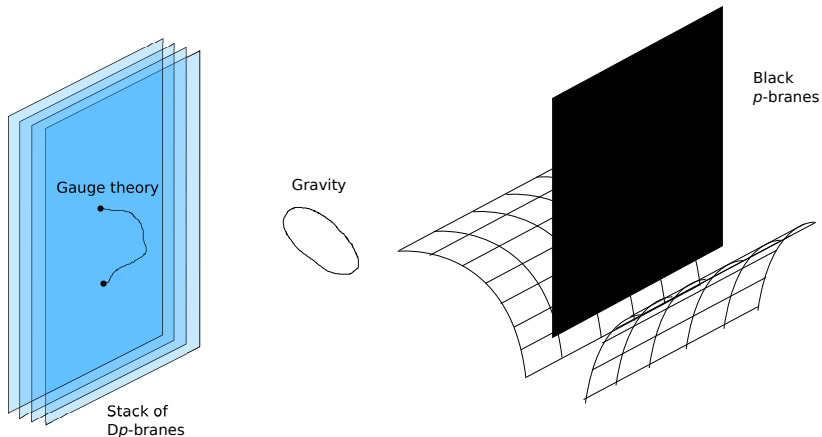
$$ds^2 = \frac{\eta_{\mu\nu} dx^\mu dx^\nu}{\sqrt{1 + (L/r)^4}} + \sqrt{1 + (L/r)^4} \delta_{ij} dx^i dx^j$$

An observer at infinity sees two different low-energy modes :

- supergravity modes in flat spacetime;
- string excitations in the throat region which is curved.

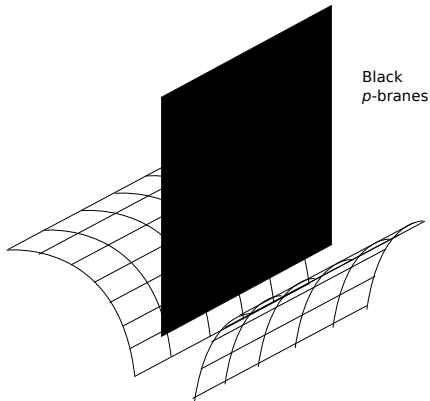
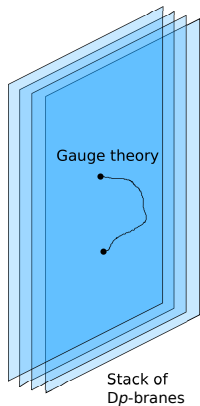


AdS/CFT correspondence



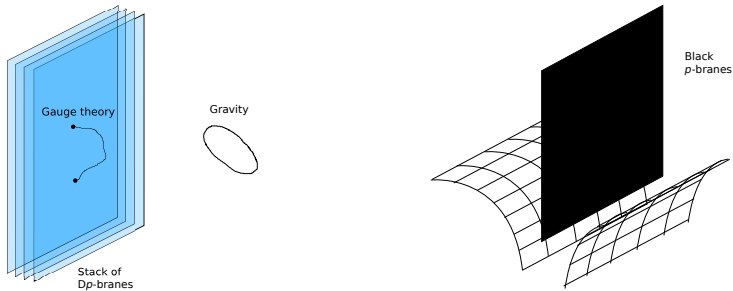
$$\left[\text{Gauge Theory} \right] \times \left[\text{Gravity in Flat space} \right] = \left[\text{String Theory on } AdS_5 \times S^5 \right] \times \left[\text{Gravity in Flat space} \right]$$

AdS/CFT correspondence



$$\left[\begin{array}{l} \text{Gauge} \\ \text{Theory} \end{array} \right] = \left[\begin{array}{l} \text{String Theory} \\ \text{on } AdS_5 \times S^5 \end{array} \right]$$

AdS/CFT correspondence

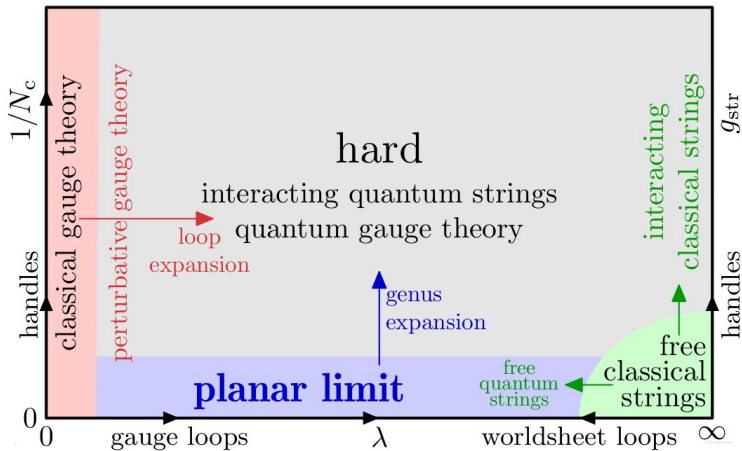


$\mathcal{N} = 4$ Super Yang-Mills theory in 4 dimensions with gauge group $SU(N)$ and coupling constant g_{YM}

Type IIB string theory with string length l_s and coupling constant g_s on $AdS_5 \times S^5$ with radius of curvature L and N units of $F_{(5)}$ flux on S^5

provided we identify the parameters as $g_{YM}^2 = 2\pi g_s$ and $2g_{YM}^2 N = (L/l_s)^4$

AdS/CFT correspondence



(Beisert et al. [1012.3982])

Thank you for your attention !

Questions ?